

# Non-Contact Imaging of Dielectric Constant with a Near-Field Scanning Microwave Microscope

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## SUMMARY

We describe a non-contact technique for imaging dielectric constant using a resonant near-field scanning microwave microscope. By measuring the shift in the system's resonant frequency as we scan over an insulating sample, we obtain quantitative images of dielectric variations. We scanned seven samples with dielectric constants  $\epsilon_r$  ranging from 1 to 230, using a 480  $\mu\text{m}$  diameter probe at a height of 100  $\mu\text{m}$  and a frequency of 9.08 GHz. The technique achieves an accuracy of about 25% for  $\epsilon_r=230$  and less than 2% for  $\epsilon_r=2.1$ , limited mainly by variations in the probe-sample separation.

## INTRODUCTION

The ability to image variations in relative permittivity or dielectric constant  $\epsilon_r$  is useful for a variety of applications. For example, in thin-film microelectronics, testing for variations in dielectric constant can be used for quality control or to develop better growth techniques. As another example, knowledge of the dielectric constant at microwave frequencies is of great importance for the design of broadband circuits. Nowadays numerous techniques exist for measuring the dielectric constant and loss tangent of insulating materials at microwave frequencies [1-6]. However, most of these techniques involve working in contact and require thick homogeneous samples [1,2,4]. In this article, we report on the use of a resonant near-field scanning microwave microscope for the non-contact imaging of dielectric samples. Our approach offers a fast, simple, broadband method to image dielectrics using readily available microwave components.

Our resonant near-field scanning microwave microscope consists of a 1 m long coaxial transmission line which is capacitively coupled to a microwave source at one end and terminated by an open-ended coaxial probe at the other end. This arrangement creates a resonant circuit in which the resonant frequency  $f_0$  and quality factor  $Q$  are modified when a sample approaches the open end of the probe (see inset in Fig. 1). By using a frequency-following feedback circuit we keep the microscope

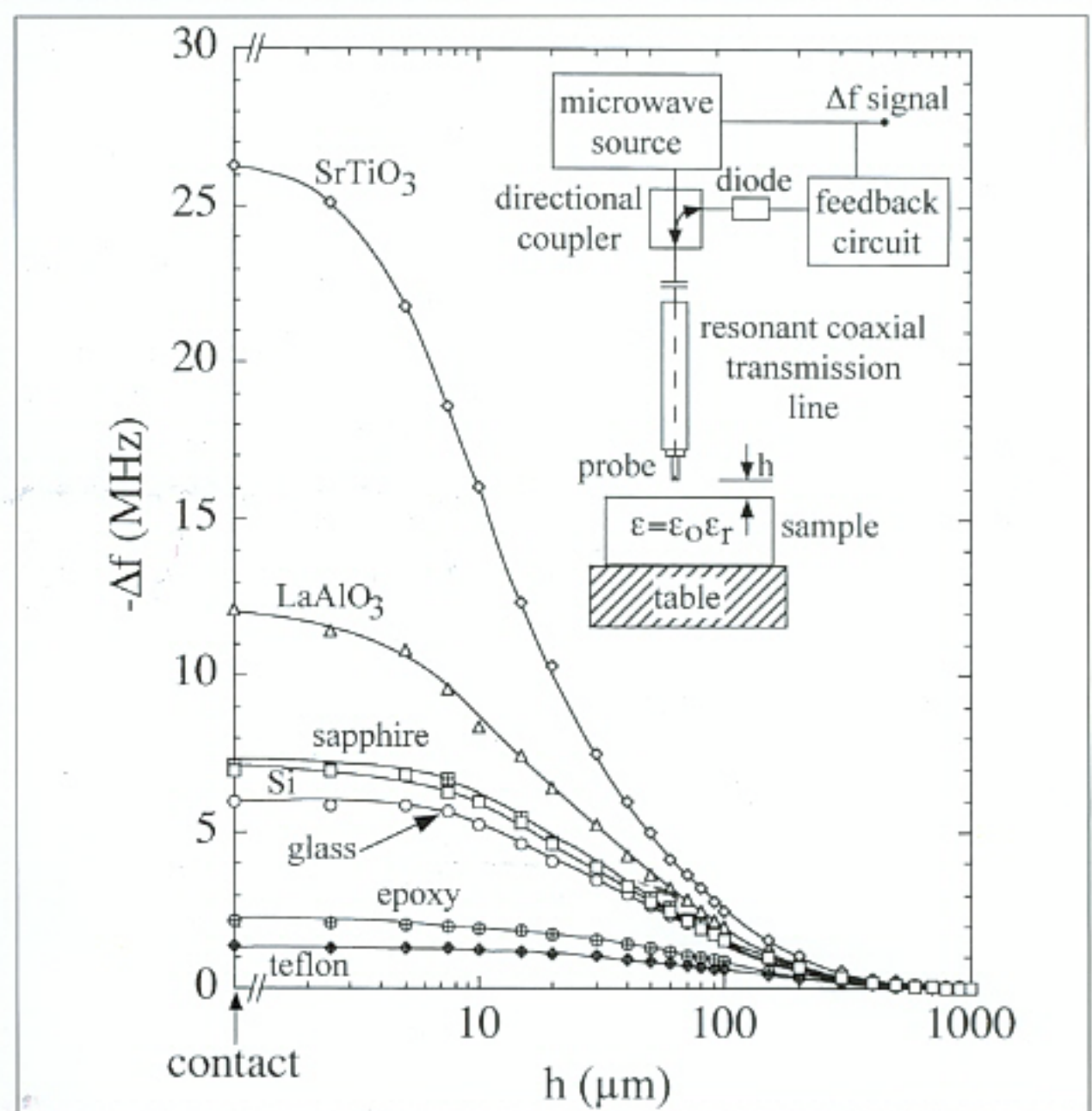


Figure 1: Frequency shift versus height for several dielectric samples measured at  $f_0=9.08$  GHz with a 480  $\mu\text{m}$  probe. Solid lines are guides to the eye. Inset shows schematic of the microwave microscope

source locked on resonance [7]. We measure the shift of the system's resonant frequency  $\Delta f$  as we scan the sample under the probe. The variations in  $\Delta f$  are directly related to spatial variations in dielectric constant in the sample. In addition, however, topographic changes will also give rise to changes in  $\Delta f$  [8].

## CALIBRATION

To calibrate the system, we constructed a test sample by placing six pieces of different dielectric material into the bottom of a square plastic mould and pouring epoxy into the mould.

In addition, silicone adhesive was used to hold each piece down. After the epoxy cured, the test sample was removed from the mould, polished, and positioned on the XY table. The materials embedded in the epoxy were silicon, glass microscope slide, SrTiO<sub>3</sub>, Teflon, sapphire, and LaAlO<sub>3</sub>. All six pieces were approximately 500  $\mu\text{m}$  thick and about 6 mm x 8 mm in size. The overall thickness of the test sample was 6 mm.

We measured the frequency shift  $\Delta f$  versus height  $h$  above the six pieces, which have dielectric constants ranging from 2.1 to about



230. We also tested the epoxy which has an unknown dielectric constant. Each piece, as well as the probe, was flat and smooth on the scale of  $5\ \mu\text{m}$  as judged by an optical microscope. For these measurements, we used a probe with a  $480\ \mu\text{m}$  center conductor diameter and a source frequency of  $9.08\ \text{GHz}$ . For each scan, the probe was first brought in contact with a dielectric and the frequency shift  $\Delta f$  was recorded as the height was systematically increased. The results are plotted in Fig. 1. Samples with the largest dielectric constant produced the largest frequency shift, as expected. The largest shift we observed was  $-26.2\ \text{MHz}$ , when the probe was in contact with a  $\text{SrTiO}_3$  sample with  $\epsilon_r=230$  [9]. The smallest shift we found was  $-1.2\ \text{MHz}$  when the probe was in contact with a Teflon sample with  $\epsilon_r=2.1$  [9]. As can be seen from Fig. 1, the frequency shift is essentially zero above  $1\ \text{mm}$  and saturates when the probe-sample distance is smaller than a few microns.

We used the above information to construct an empirical calibration curve that directly relates the frequency shift to the dielectric constant. In order to construct the calibration curve we took the difference between the frequency shift at two different heights  $h_1$  and  $h_2$ , i.e.  $f_d = \Delta f(h_2) - \Delta f(h_1)$ , where  $h_2$  is far away ( $h_2 > 1000\ \mu\text{m}$ ). By taking the difference, we eliminated the effect of drift in the microwave source frequency. Proceeding this way for the test samples, we constructed two calibration curves of  $f_d$  versus  $\epsilon_r$  (see Fig. 2), one curve for  $h_1=10\ \mu\text{m}$  and  $h_2=1.1\ \text{mm}$  and the other for  $h_1=100\ \mu\text{m}$  and  $h_2=1.1\ \mu\text{m}$ . We then set the parameters in each calibration curve with an empirical function (solid lines in Fig. 2), allowing us to easily transform any measured frequency shift to a dielectric constant. From these curves we can see that we can enhance the sensitivity to the dielectric constant considerably by using a small probe height. On the other hand, at closer probe-sample separations the influence of topographic features will be enhanced.

## IMAGING RESULTS

To test the dielectric imaging capabilities of our system, we next scanned a single sample of  $\text{LaAlO}_3$  which had an  $8 \times 5\ \text{mm}$  triangular shape and a thickness of  $510\ \mu\text{m}$ . We placed the sample directly on the metal scanning table and recorded the frequency shift as a function of position. The data was taken at  $9.08\ \text{GHz}$  using the  $480\ \mu\text{m}$  probe at heights of  $100\ \mu\text{m}$  and  $1.1\ \text{mm}$ . We subtracted the two data sets and used our  $100\ \mu\text{m}$  calibration curve to transform the resulting frequency shift image into a dielectric constant image. Figure 3 shows the resulting contour image of dielectric constant versus position. The dielectric constant varies from about 20 to 25 over the sample and equals 1 when the probe is away from the sample. For comparison, the reported value of relative permittivity for  $\text{LaAlO}_3$  at room temperature is  $\epsilon_r=23.9$  at  $18\ \text{GHz}$  [10]. In this image, the main variation in  $\epsilon_r$  over the sample is due to a slight tilt in the sample surface of about  $20\ \mu\text{m}$ . The edges of

Table 1:

Experimental measurements of dielectric constant found at the centre of each material from Fig. 4 compared to literature values.

Material	Experimental value	Literature value	Frequency (GHz)	References
Silicon	12	11.7	100	11
Glass	12	6.7	$10^{-3}$	12
$\text{SrTiO}_3$	180	230	0.1	9
$\text{LaAlO}_3$	20	23.9	18	10
Sapphire (ceramic)	20	10.0	100	11
Teflon	2.1	2.1	10	9

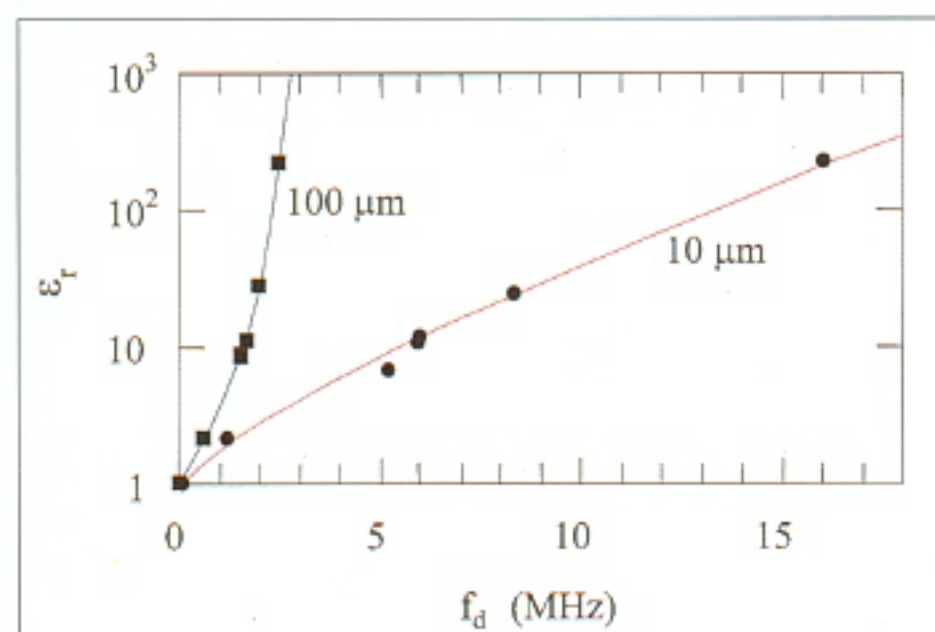


Figure 2: Calibration curves of dielectric constant  $\epsilon_r$  versus frequency shift  $f_d$  using a  $480\ \mu\text{m}$  probe at  $9.08\ \text{GHz}$  for closest approach heights of  $100\ \mu\text{m}$  and  $10\ \mu\text{m}$ .

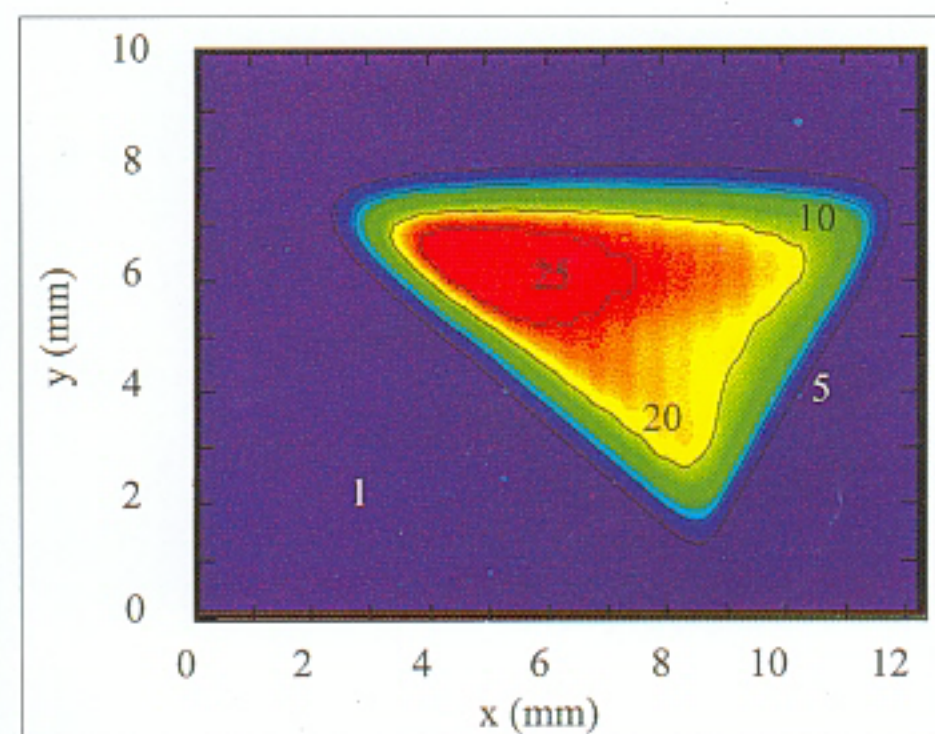


Figure 3: Dielectric constant of  $\text{LaAlO}_3$  sample imaged at  $100\ \mu\text{m}$  and  $9.08\ \text{GHz}$  using a  $480\ \mu\text{m}$  diameter probe. Contour lines are labelled with dielectric constant.

the sample show a smaller value of  $\epsilon_r$  due to averaging over the inner conductor of the probe. The width of the affected region is in good agreement with the expected spatial resolution of about  $500\ \mu\text{m}$ .

We next recorded a frequency shift image of the six-piece test sample using a probe-sample separation of  $100\ \mu\text{m}$  and the  $480\ \mu\text{m}$  diameter probe at  $9.08\ \text{GHz}$ . As before, using the  $\epsilon_r(f_d)$  calibration, we then transformed the frequency shift image into a dielectric constant image (see Fig. 4a). The darker regions in Fig. 4a indicate a higher dielectric constant (larger frequency shift) and the lighter regions indicate a smaller dielectric constant (smaller fre-

quency shift). Figure 4b shows the corresponding surface plot representation. Note that the z-axis in Fig. 4b uses a logarithmic scale to allow us to show the large range of dielectric constants present in the sample. As expected, the largest dielectric constant materials ( $\text{SrTiO}_3$  and  $\text{LaAlO}_3$ ) are the highest surfaces and the smallest dielectric constant materials (Teflon and vacuum) are the lowest. Further, notice that the Teflon sample forms a depression in Fig. 4b, indicating that the dielectric constant of Teflon is lower than the surrounding epoxy. We also note that voids in the epoxy can easily be seen as irregularly shaped low-dielectric regions (white regions



in Fig. 4a) in the epoxy. Table 1 summarizes the dielectric constants found from Fig. 4 for the six test materials.

From Fig. 4b, it is apparent that there is some noise in our images of dielectric constant. In our system, the predominant sources of random errors are noise in our recording electronics and fluctuations in the source frequency. To establish the precision to which  $\epsilon_r$  can be determined, we determined the standard deviation in  $\epsilon_r$  over a small region near the center of the  $\text{LaAlO}_3$  sample in Fig. 3; we found  $\Delta\epsilon_r=0.06$  for a sampling time of 30 ms. Similarly, over the Teflon in Fig. 4a the standard deviation in  $\epsilon_r$  was  $7 \times 10^{-4}$  for a sampling time of 30 ms. From Fig. 4 we can also estimate the absolute accuracy of our technique. For  $\epsilon_r=230$  ( $\text{SrTiO}_3$ ) the accuracy is about 25% while for  $\epsilon_r=2.1$  (Teflon) the accuracy is better than 2%. The main source of these errors is topographic variations. In our test sample, even after polishing, there are small height variations of about  $30 \mu\text{m}$  (e.g. between the sapphire and  $\text{SrTiO}_3$  in Fig. 4) between the different dielectrics. Such height variations cause an additional frequency shift [8], resulting in an error in the measured dielectric constant. In this regard, an accurate measurement of the dielectric constant of a single flat dielectric sample is considerably easier. However, the image of the composite sample clearly demonstrates the strength and sensitivity of this non-contact technique to measure variations in  $\epsilon_r$ .

## CONCLUSIONS

In summary, we have used an open-ended, resonant, near-field scanning microwave microscope to obtain quantitative images of the dielectric constant of bulk materials. Our system allows for fast, reliable, non-contact imaging of variations in the dielectric constant of flat samples, provided that the height of the probe above the sample is accurately controlled.

## ACKNOWLEDGEMENTS

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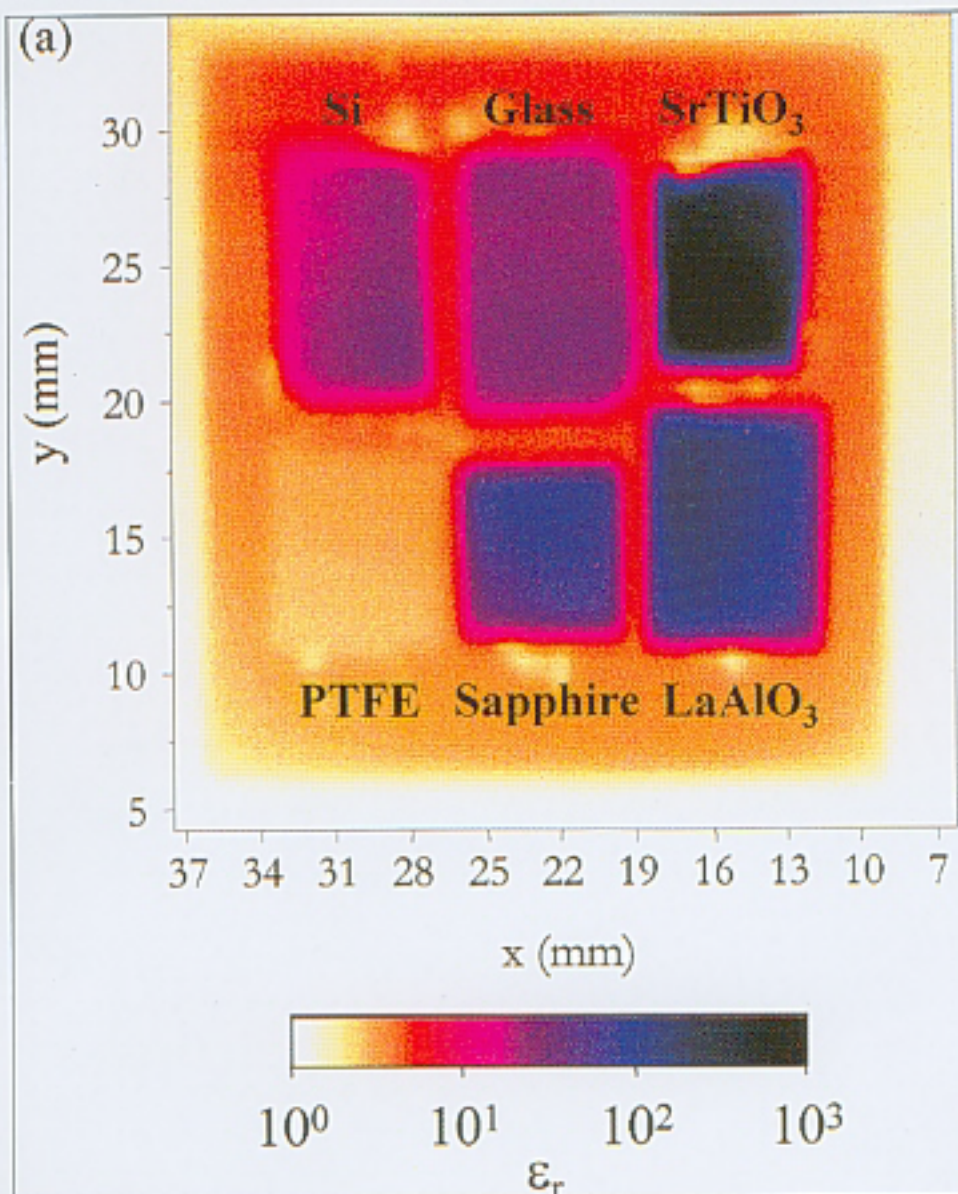


Figure 4a

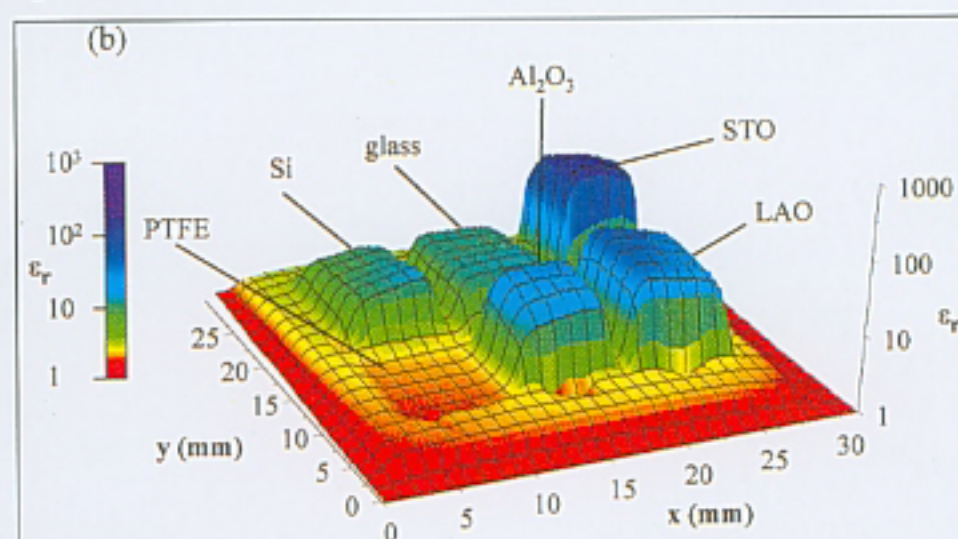


Figure 4b

Figure 4:

a: Dielectric constant image of a test sample taken with a  $480 \mu\text{m}$  probe at 9.08 GHz and  $100 \mu\text{m}$  above the sample. The materials used in the sample are, along the top row from left to right: silicon, glass,  $\text{SrTiO}_3$ ; and along the bottom row from left to right: Teflon, sapphire,  $\text{LaAlO}_3$ .

b: Surface plot of the above sample on a logarithmic z-scale.